

that the relation between  $\bar{P}_u$  and  $K$  is essentially independent of  $T_p$ . It appears that  $K_n^*$  shifts to lower values because of the temperature displacement of the stable pressure. Thus, although firings at high temperatures were stable, the motor became inherently unstable at a sufficiently low value of  $T_p$ .

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FEBRUARY 1964

AIAA JOURNAL

VOL 2, NO 2

## Comparison of Commercial, Spherical Powder, and Wire Bundle Tungsten Ionizers

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Cesium ionization on different porous tungsten structures was measured in terms of neutral fraction, critical temperature, and flow rate in the  $10^{-6}$ -torr range. These structures included  $\frac{3}{16}$ -in commercial sintered ionizers with about  $2\text{-}\mu$  average pore diameter, spherical powder ionizers of  $6\text{-}\mu$  effective pore diameter, and 12- and  $6\text{-}\mu$ -diam wire bundle ionizers. Conductance was measured at room temperature with nitrogen and at operating temperature with cesium vapor. Ion current and neutral fraction were plotted as  $\log j$  or  $\log \alpha$  vs  $1/T$ , which allows easier extrapolation of the critical temperatures. The directly heated ionizers could be operated up to  $1500^\circ\text{C}$ . In general, the performance of ionizers improved with time of operation. During this conditioning time, the critical temperature and neutral fraction decreased. The best critical temperatures for all tungsten ionizers in the  $10^{-6}$ -torr residual gas pressure at current densities between 1 and  $30\text{ ma/cm}^2$  were within  $100^\circ\text{K}$  of Langmuir's data for solid tungsten. The critical temperature had a considerable spread. Neutral fractions for commercial sintered, spherical powder, and fine wire bundle ionizers ranged from 0.1 to 15% but were considerably higher for coarse wire bundle ionizers. The angular distribution measurements of neutral fraction were used to select the best angular position of a neutral efflux detector for neutral fraction measurements and to test the validity of the assumption that the relative angular distribution remains independent of whether a beam is being extracted.

### I Review of Surface Ionization on Porous Tungsten

THE first practical cesium ionizers in the early ion sources, surface ionization detectors, and even ion propulsion devices used solid tungsten filaments or ribbons. Surface ionization of the alkali metals on solid refractory metal surfaces

has been reviewed recently by Zandberg and Ionov.<sup>1</sup> A brief review of surface ionization on other solid materials was given previously.<sup>2</sup>

Studies of porous ionizers started with the application of commercial porous tungsten in ion engines in late 1958. At the same time, theoretical studies of surface ionization on idealized porous structures determined that the desirable pore size should be less than  $1\mu$ .<sup>3-5</sup>

More basic experiments with porous ionizers were performed later. The first of these, by Staviskii and Lebedev,<sup>6</sup> used a relatively large (10-mm diam, 2-mm thick) porous tungsten disk of 70% of theoretical density, made of a very fine  $1\text{-}\mu$  grain size.

The ion current-ionizer temperature curves were similar to those for solid tungsten. Following the definitions of dif

Presented as Preprint 63016 at the AIAA Electric Propulsion Conference, Colorado Springs, Colo., March 11–13, 1963; revision received October 7, 1963. This work has been supported by NASA Contracts NAS8-2547 and NAS8-1537 and Air Force Contract AF 33(616)-6958. Spherical powder and wire bundle ionizers were developed by H. Todd and M. LaChance.

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ferent regions of surface ionization, the ion current in the low coverage region was constant, suggesting near 100% ionization. The critical temperature  $T_{CI}$  was  $80^\circ$  above Langmuir's data for solid tungsten at  $10 \text{ ma/cm}^2$  and  $50^\circ$  higher at  $0.25 \text{ ma/cm}^2$ . The pressure was not specified, but an apparatus diagram suggests an untrapped oil diffusion pump system with operation in the  $10^{-6}$ -torr range.

The Faraday cup collector was also the accelerator electrode with its aperture about 7 mm from the ionizer. Thus, sputtering of cesium and the unidentified collector material could reach the ionizer, since there was no baffling to minimize it.

More extensive studies of porous tungsten ionizers at pressures below  $10^{-7}$  torr were reported by Husmann. Operation in a bakable, mercury- and ion-pumped glass system in the  $10^{-8}$ -torr region after a short flash of the ionizer at  $1700^\circ\text{C}$  produced a "clean" surface.

The critical temperatures for fine porous disks ( $2\text{-}\mu$  pore diam) at  $10 \text{ ma/cm}^2$  were only  $40^\circ$  to  $80^\circ\text{C}$  above Langmuir's data for solid tungsten. For coarse disks or fine oxidized disks, the critical temperatures were about  $150^\circ\text{C}$  higher than Langmuir's. The neutral fractions for clean, fine disks ranged from 0.5–2% at  $1 \text{ ma/cm}^2$  to 7–9% at  $10 \text{ ma/cm}^2$ . For oxidized disks, the neutral fractions were  $\frac{1}{3}$  lower. Oxidation was provided by admitting oxygen to raise the total pressure to  $10^{-6}$ – $10^{-5}$  torr.

Emissivity of sintered tungsten was measured as a function of pore size. The emissivity increased with temperature from  $100^\circ$  to  $1200^\circ\text{C}$  with a slope similar to that of solid tungsten. The emissivity of disks made of 0.9- to  $18\text{-}\mu$  powder at  $1200^\circ\text{C}$  ranged from 0.25 to 0.55 independently of powder size. Mass spectroscopic analysis of the residual gas at  $10^{-7}$  torr indicated the presence of  $\text{H}_2$ ,  $\text{N}_2$ , and  $\text{H}_2\text{O}$  but no  $\text{O}_2$  or  $\text{CO}_2$ . The sensitivity was sufficient to give an indication at  $2 \times 10^{-9}$  torr.

The first attempt to study surface ionization as a function of ionizer grain size (or structure) gave ionization efficiencies of 94 and 97% for ionizers made of 5- and  $8\text{-}\mu$  powder, respectively. Since no ion current density, critical temperature, or emissivity were given, this work could not be compared with that of other investigators.<sup>8</sup> The atom and ion emission from a much larger 1-in.-diam porous tungsten ionizer was studied by Shelton<sup>9</sup> using a tantalum ribbon neutral sensor. At  $10^{-6}$  torr and ion current densities of about  $0.5 \text{ ma/cm}^2$ , the neutral fraction ranged from 0.2 to 3%. The tungsten was often oxygenated, and some cases of poisoning were observed. Some data by these investigators are included in the graphs of the final section.

## 2 Ionizer Performance Criteria

The important ionizer performance criteria<sup>10</sup> are ionization efficiency, ion generation energy efficiency, and lifetime. Since the ionization efficiency is usually very close to 100%, more accuracy is obtained by measuring the neutral fraction, defined as the ratio of neutral cesium efflux to total cesium efflux. The ion generation energy efficiency is a measure of thermal losses due to radiation from the ionizing surface, radiation and conduction losses from the supporting structure, and leads. For evaluation of the ionizer material, only the thermal radiation from the ionizing surface is of importance. However, in a small single-button device other losses predominate. The radiation loss from the active area is best obtained by measurement of the minimum operating temperature and the total emissivity. The minimum operating temperature is equal to the so-called critical temperature. For porous tungsten ionizers in  $10^{-6}$ -torr vacuum, most good ionizers have a rather sharp critical temperature and a very small difference between the lower and upper critical temperature (hysteresis) determined from the ion current-ionizer temperature curve. For many ionizers, ion current starts a sharp decline at the same time as the neutral fraction starts to

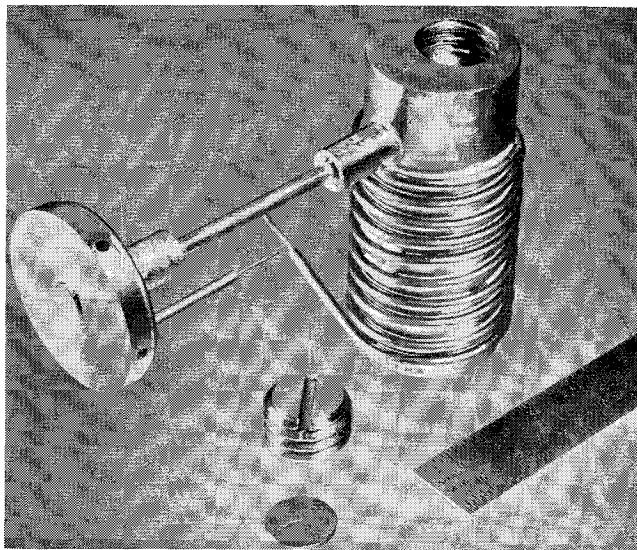


Fig 1 Ionizer reservoir assembly

rise. Some ionizers, however, showed a minimum in neutral fraction at a temperature higher than the critical temperature for the ion current. The ion current critical temperature  $T_{CI}$  is most conveniently defined as that temperature at which the current has decreased by 5% from its maximum value. This is discussed in more detail below.

The total emissivity can be measured in a separate experiment using the same material. It is assumed here that the differences due to emissivity are smaller than the effect of the spread in the critical temperatures and that they would favor high-emissivity ionizers.

Some indication concerning the life of the ionizers can be obtained by observing time trends in their gas and cesium permeability, neutral fractions, and critical temperatures. Only limited data were obtained because of the short duration of most tests. Many of these tungsten ionizers were operated for over 100 hr between  $1400^\circ$  and  $1500^\circ\text{C}$  without any decrease of permeability.

## 3 Ionizer Test Apparatus

Nearly all ionizers tested were disks of  $\frac{3}{16}$ -in. diam  $\times$  0.040-in. thick. They were brazed into a molybdenum holder, as shown in Fig 1, using an Mo-C-B braze. This ionizer assembly was mounted into a single-aperture ion accelerator using a  $45^\circ$  beam-forming electrode and planar accel-decel electrodes, as shown in Figs 2 and 3.

A few grams of cesium were loaded into the reservoir using regular dry box techniques. Upon removal of the ionizer assembly from the dry box, the cesium in the reservoir was under 1 atm pressure of nitrogen. The low permeability of the porous ionizer was the only barrier against the infusion of air while the assembly was mounted into the accelerator and transferred into the vacuum test chamber.

The ionizer-accelerator assembly was placed in a  $1 \times 3$ -ft vacuum tank that was equipped with either a stationary or a movable neutral detector, a liquid-nitrogen-cooled ion collector, and a liquid-nitrogen-cooled liner. Operating pressure was in the  $10^{-6}$ -torr range. A schematic view of the neutral eye is also shown in Fig 2 and the ion collector in Fig 4. The ionizer temperature was measured by a W5%Re-W26%Re thermocouple in a  $\frac{1}{16}$ -in.-o.d. tantalum sheath. The hot junction was in a deep hole in the ionizer holder and very near to the ionizer disk. The other end of the thermocouple passed through the vacuum wall without interruption by means of a gasketed feedthrough.

The cesium reservoir temperature was monitored by a bare iron-constantan thermocouple near the exit of the massive nickel reservoir.

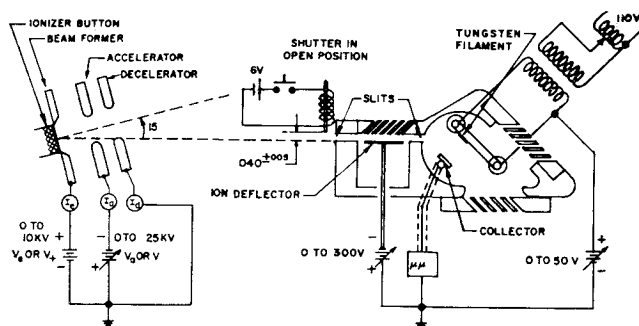


Fig 2 Schematic and circuit diagram of ion accelerator and neutral detector

#### 4 Ionizer Test Procedure

The gas conductance of most ionizers was measured before and after brazing into the holder and again after the ionizer tests. The ionizer and reservoir heaters were turned on at the same time as all instrumentation and high voltages. The neutral eye filament was baked out at 1600°C until the background reading with the shutter closed was a few nanoamps. The operating temperature was between 1100° and 1300°C, depending upon atom incidence rates. The ionizer was heated at about 1300°C.

The ion collector current was monitored to establish the proper cesium-reservoir temperature for a reasonable ionizer-current density, say, 5 ma/cm<sup>2</sup>. With the reservoir temperature and cesium flow stabilized, the accelerator/ionizer voltage ratio for proper focusing was checked visually by observing the ion beam impact area at the collector or by observing the beam at 10<sup>-5</sup> torr of argon. Usually the accelerator voltage was twice the value of the ionizer voltage, and the beam was 1-2 in. in diameter at the collector.

With proper focusing, the current-voltage characteristic was taken to insure emission-limited operation. This was done for each new current density and is very important at the high current densities when operation is very near to being space-charge limited. The movable neutral detector was usually located at 15° to the ion beam axis and 6 in. from the ionizer. The movable detector had 0.020-in.-wide slits. The stationary detector had 0.040-in. slits and was only 3 in. from the ionizer. If the neutral fraction was below 10%, indicating that the ionizer had been properly conditioned, the angular dependence of  $\alpha$  was checked to see that it was fairly constant and that the probe was located at the proper angle for reliable results. Operation too close to the beam or

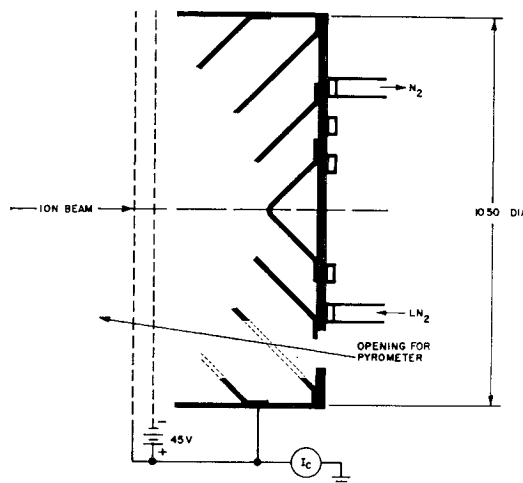


Fig 4 Conical baffle current collector with electron impression

too much in the shadow of the electrodes gave erroneous results.

The neutral fraction was obtained by measuring the neutral detector current with accelerating voltage on and off, subtracting background readings obtained by operating the shutter. After the proper probe angle was checked, the ion current and neutral fraction were measured as a function of ionizer temperature for a series of current densities. The dependence of cesium permeability upon the ionizer temperature was also obtained, since it is related to the ion current density  $j$ . The  $j$  vs  $T$  curves were taken repeatedly to observe reproducibility, hysteresis, and time trends.

After each test, the ionizer was inspected and its nitrogen conductance checked. Three types of tungsten ionizer structures were tested: 1) commercial sintered ionizers designated by the symbol *SW*, 2) spherical powder ionizers designated *SP*, and 3) wire bundle ionizers designated *WB*.

#### 5 Ionizer Nitrogen Conductance Permeability and Braze Infiltration

Nitrogen conductance measurements using electronic timer equipment provided a fast and simple method to check and sort large numbers of unmounted porous ionizers. The actual measurement determined the time for the pressure in a constant volume to drop about 1% of the pressure difference across the ionizer. The inverse of this time is proportional to the conductance of the sample. Conductance depends upon both the area and the thickness. Both can be easily measured for unmounted ionizers and used to calculate permeability. The permeability value is independent of sample area and thickness and depends only upon the material. The distribution of permeability for a large number of ionizers is indicative of the ionizer material uniformity. The gas permeability is a much more sensitive indicator of material quality than apparent material density. Change in permeability with time at high temperature provides an indication of structural instability. For ionizers brazed into a holder, the open area is no longer definite due to braze infiltration, and it is impossible to calculate gas permeability. One can assume that the permeability

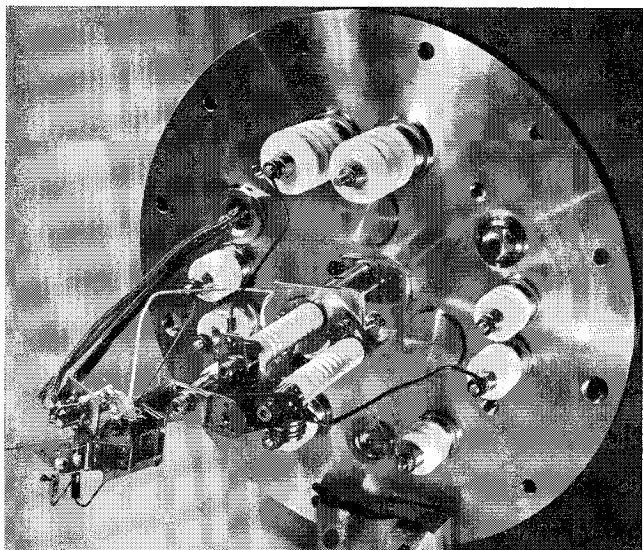


Fig 3 Ionizer testing equipment including neutral detector

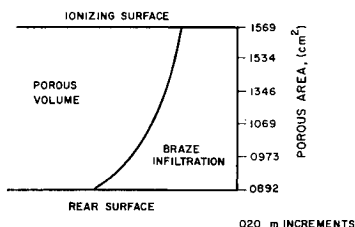


Fig 5 Braze penetration into spherical powder ionizer as measured from a cross section metallograph, SP2,  $A = 0.179 \text{ cm}^2$

**Table 1 True and apparent permeability for spherical powder ionizers**

Ionizer identity	True permeability	Apparent permeability after brazing $K_1$ , g/(cm-torr-sec)	Apparent permeability after test $K_2$ , g/(cm torr-sec)	$(K - K_1)/K = (A - A_1)/A$	Cesium flow $j$ , ma/cm <sup>2</sup> at 300°C, $A_1/0.157$ cm <sup>2</sup>
	$K = \frac{\Delta m}{\Delta p} \frac{\Delta x}{A} \frac{1}{\Delta t}$ , g/(cm-torr-sec), $\Delta x/A = 0.5705$				
SP1	$1.42 \times 10^{-6}$	$0.37 \times 10^{-6}$	$0.437 \times 10^{-6}$	0.74	3.50
SP2	0.92	0.493	0.632	0.462	4.47
SP3	1.17	0.628	0.624	0.464	3.78
SP4	1.17	0.553	0.488	0.528	1.0
SP5	0.92	1.21		(-0.315)?	
Average	1.12	0.650	0.554	0.548	

of an unmounted ionizer does not change during brazing, if the brazing temperature is considerably lower than the sintering temperature.

The nitrogen conductance was checked before performance testing. These tests are also useful for the detection of leaks after brazing, cracks upon testing, braze penetration, or changes of permeability. Care must be taken to remove all cesium, its oxides, and cleaning fluid from the ionizer after testing, since these can effectively seal the ionizer pores. Most ionizers were baked out in vacuum after their performance testing.

The conductance was measured before and after brazing the ionizer into a holder. The permeability was calculated from the unmounted ionizer measurement. An apparent permeability was calculated for the mounted ionizer, using the same area and thickness value. The ratio of the true and apparent permeability then gives the ratio of initial to effective post brazing area. Table 1 shows true permeability, apparent permeability before and after performance testing, and the closed area fraction and cesium permeability represented by  $j$  at 300°C for the SP ionizers. From this table the fraction of ionizer area closed by brazing varies from 46 to 74% with an average 55% closure. The ionizer holder produces a braze infiltration pattern, as shown in Fig. 5. From Fig. 5 and the column of porous area, one can see the extensive braze infiltration at the rear surface. The gas conductance should be determined by the average cross-sectional area near the center of the ionizer. The braze closure varies from 12% at the front surface to 50% at the rear surface, whereas the gas conductance closure was 46%. Hence, the gas conductance is determined more by the minimum open cross section than by a geometric average.

The apparent nitrogen permeability of the average SP ionizer is two to three times that of the average SW ionizers. The WB ionizers are about 10 times more permeable.

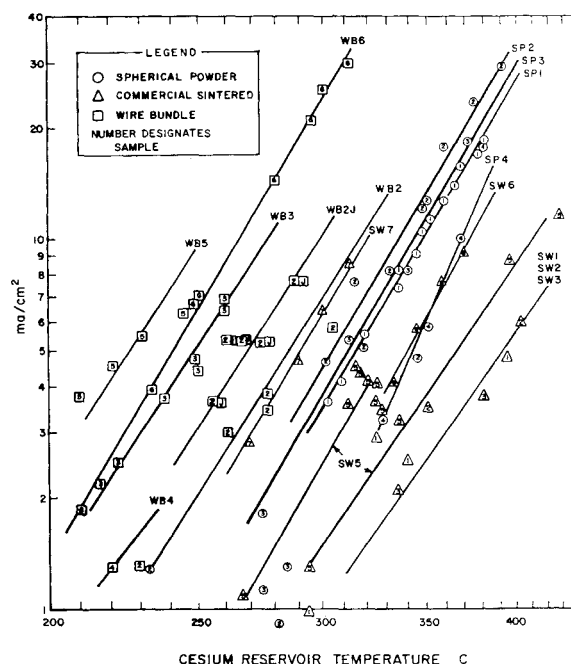
The commercial sintered tungsten ionizers (designated by the symbol SW) were selected from a large number of buttons to have a range of permeabilities convenient for a reasonable cesium reservoir temperature and a fine and stable pore structure. The average spherical powder ionizers (designated by the symbol SP) and the wire bundle ionizers (designated by WB) were made by H. Todd and M. LaChance of Electro-Optical Systems, Inc., under a NASA contract.

## 6 Cesium Permeability and Its Change with Temperature

A measure of cesium conductance can be obtained in two ways: 1) by using the cesium atom probe, and 2) by measuring cesium ion current. The neutral probe current was measured for two SW and two WB ionizers used for the study of angular distribution of the cesium effusion pattern. The probe current on the ionizer axis was assumed to be directly proportional to the total cesium flow in particles per second. Data showed that the cesium conductance of WB ionizers was about 50 times that of SW ionizers. Each group has a spread of factors up to 10.

The second type of cesium conductance estimate is the ion current density in Fig. 6, which does not include neutrals. For SW and SP ionizers, the actual cesium conductance in equivalent milliamperes per square centimeter was 1-15% higher and for WB about 5-50%. Note also that higher current densities required higher ionizer temperatures. The cesium conductance is slightly dependent upon the ionizer temperature, as seen in Fig. 7. The ionizer was not held at constant temperature in Fig. 6. The ion current measurement gives the value of cesium conductance which is converted to an apparent cesium permeability measurement by normalizing to equal areas for each of the ionizer groups. The best estimates of the active areas are 0.157 cm<sup>2</sup> for SW and SP ionizers, except for SW1, which was 0.220 cm<sup>2</sup>, and 0.141 cm<sup>2</sup> for WB ionizers. The SP ionizers have apparent cesium permeabilities three times higher and the WB about 10 to 20 times higher than the SW ionizers.

The change of cesium conductance with temperature can be measured with the atom probe, assuming that the angular distribution does not change with ionizer temperature. Figure 8 shows the decrease of ionizer permeability with temperature as measured by the neutral probe on the geometric axis. In this case there is no voltage applied and no surface ionization, and there is a high surface coverage. The equivalent ion current density indicated for each curve gives the average ion current that would be generated at this flow rate (or cesium temperature) with 100% ionization. The temperature coefficients of conductance shown were computed at  $T = 1200^\circ\text{C}$ .



**Fig. 6 Cesium current density as a cesium permeability indicator for the various tungsten ionizers**

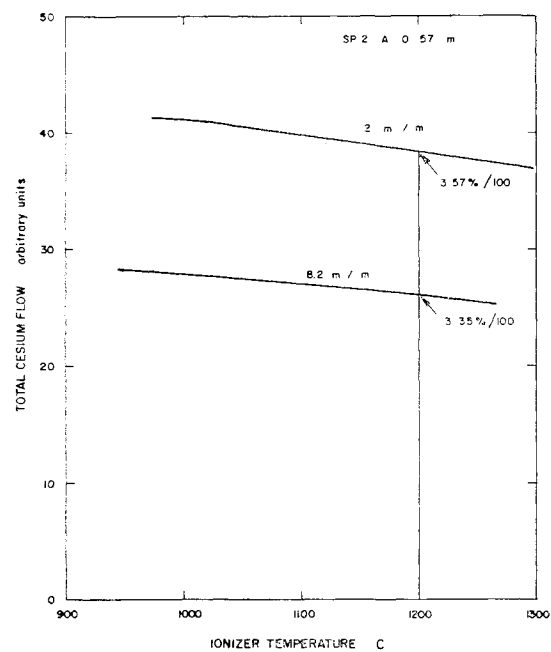


Fig 7 Change of cesium permeability with ionizer temperature at high surface coverage

7 Active Ionizer Area and Ion Current Density Measurement

Measurement of the active ion-emission area is needed for the calculation of ion current density. The active area can be estimated 1) by microscopic observation of ionizer surface for a color change around the edges; 2) by comparison of nitrogen conductance before and after brazing into the holder; 3) by obtaining a sputtering pattern in a parallel plane arrangement; 4) by low magnification ion microscopy; and 5) by a photomicrograph of the ionizer cross section, which destroys the ionizer. The first method is the simplest but is not very reliable. The second method gives the relation between true and apparent nitrogen permeability but does not give an accurate measure of active area as discussed in the previous section.

A more reliable method is the measurement of active emission area by the sputtering pattern. The ionizer is placed into a plane parallel diode arrangement with a copper cathode and a gap of about 0.175 in. Emission-limited operation is maintained to assure parallel ion trajectories. This produces

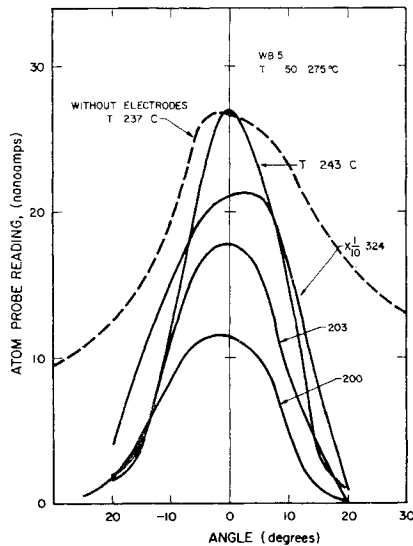


Fig 8 Angular distribution of atoms through electrodes

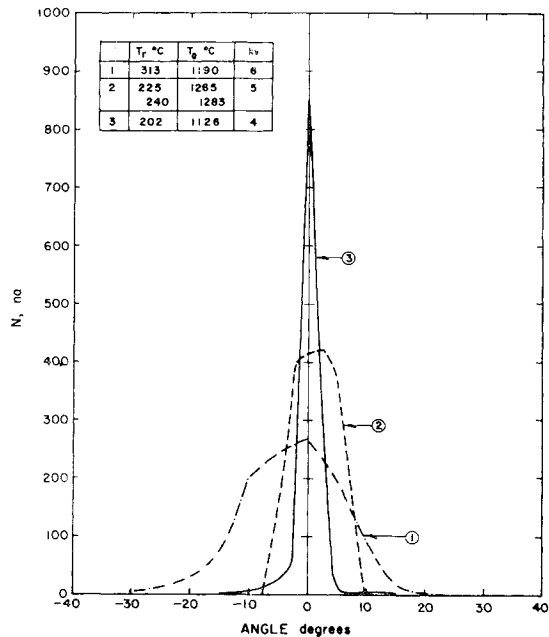


Fig 9 Ion beam profiles for a wire bundle ionizer

a well defined sputtered area that is assumed to be equal to the active ion-emission area. The sputtering pattern area for SW7 was 0.155 cm<sup>2</sup>, compared to 0.157 cm<sup>2</sup> measured by metallographic cross section for SP2. Based on these two measurements, active areas for early SW and SP ionizers were assumed to be 0.157 cm<sup>2</sup> and measured for all recent ionizers.

For the wire bundle the surface closure measured on a metallographic cross section was 15.2%, and the penetration pattern was different. The 0.141-cm<sup>2</sup> area was used for most 12-μ wire bundles. For the fine wire bundle, the sputtering pattern gave an area of 0.114 cm<sup>2</sup> or a 23% closure.

The measurement of ion current is made at a collector and involves the following assumptions:

1) The entire beam emerging from the engine is collected. This is checked by a visual observation of the target luminescence area, usually from 1 to 3 in in diameter. The collector diameter is 10 in. It can also be checked by visual observation of the beam angle in argon. A third check is provided by the I-V characteristic. When some of the beam misses the target, this fraction usually depends upon the voltage and produces an irregular I-V curve.

2) There is no secondary electron-emission current away from the collector. This is checked by increasing the secondary electron-suppressor voltage until the collector current no longer decreases.

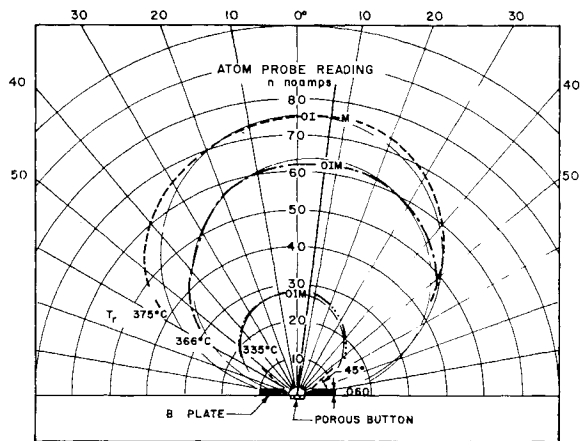


Fig 10 Angular distribution of cesium atoms evaporating from a commercial sintered porous tungsten ionizer as a function of cesium flow rate, SW3, T = 1200 C

3) There is negligible interception of the emitted ions by the accelerating structure. Ion interception at the accel electrode is small but measurable. Measurements are complicated by electron currents.

The active surface area measurement includes the following additional assumptions:

1) Emission is from the apparent geometrical area, without regard to the ratio of pore area to closed area or to any true solid area dependent on roughness and grain structure of the surface.

2) There is no emission from the brazed area or the holder. Most of the preceding factors would tend to increase the current density; hence our measurements reflect a conservative value.

A typical current-voltage characteristic is shown in Fig. 9. The ratio of accelerator  $V_-$  to ionizer  $V_+$  voltage is kept constant and equal to 2. This is equal to an accel-decel ratio  $(V_- + V_+)/V_+$  of 3. The accelerator perveance for this curve is 1.0 nperv of 6.5 nperv/cm<sup>2</sup>, using for the area the active area of 0.157 cm<sup>2</sup> corresponding to a braze infiltration of 12%.

### 8 Angular Distribution of Atoms, Neutrals, Ions, and Neutral Fraction

An important ionizer performance figure of merit is the neutral fraction. The neutral fraction is defined as the ratio of neutrals to the total efflux of cesium ions and neutrals. The term "neutrals" in this report will be used specifically for the case where the accelerating voltage is applied and an ion beam is being extracted. The term "atoms" will be used for the case where no accelerating voltage is applied and the entire efflux is atoms.

Since, with the use of a probe, a small portion of the neutral fraction rather than the total is measured, one must ascertain whether the neutral fraction is independent of the angle. For this reason the angular distributions of neutrals, atoms, and neutral fractions were studied in detail. In some of the measurements the slit widths were 0.040 in. with a notched tungsten filament. In some of the latter experiments the slit widths were reduced to 0.020 in. and the notches eliminated in order to reduce the chance of systematic angular error.

The angular distribution of cesium atoms evaporating from a commercial sintered porous tungsten ionizer without any accelerating structure is shown in Fig. 10 as a function of the flow rate. This distribution is very nearly cosine. It is typical for most measurements of this type, and it was independent of the cesium flow rate and of the ionizer temperature. The angular distribution of cesium atoms evaporating from a wire bundle showed a very definite peak, and it deviated more from the cosine law depending upon the nature of the mechanical surface treatment. Figure 11, for example, shows a highly peaked pattern obtained with an etched wire bundle. The etching has opened up the pores. Distributions for some of the other wire bundle ionizers were not as highly peaked because their surface pore openings were reduced by mechanical abrading or by compression of the button. This definitely suggests that the atoms evaporating from a wire bundle come to some degree from within the pores. A shielding effect of the electrode on the distribution of atoms is shown in Fig. 8 or from comparison with the curve for the case without electrodes. A complete study of the angular distributions of ions, atoms, neutrals, and neutral fraction through electrodes for a wire bundle WB5 is shown in Fig. 12. These curves illustrate that the neutral fraction cannot be measured with this probe inside of the ion beam, which is about 18° wide. The signal due to ions is considerably larger than that due to atoms or neutrals. The neutral detector can be used for a study of ion beam profiles illustrated in Fig. 9 on a rectangular plot. These curves illustrate the ion beam ex-

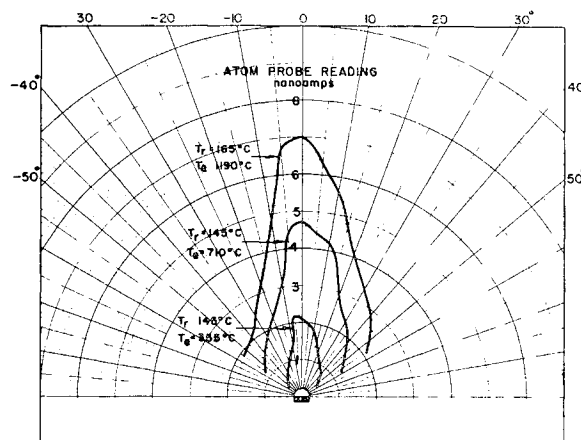


Fig. 11 Angular distribution of cesium atoms as a function of ionizer temperature, etched WB5

pansion from a total included angle of 8° to about 30° as the current density is increased from about 3 to about 12 ma/cm<sup>2</sup>, even though the voltage is also increased.

Either a single screen accelerating electrode or a set of slotted accel-decel electrodes was used to measure the angular dependence of the neutral fraction. In both cases, the neutral fraction appeared reasonably constant within an angle of about 20° just outside of the ion beam and before the signal strength fell off to a very low value. Most of the checks of the constancy of neutral fraction with angle were done in the presence of the regular electrodes. A set of curves for a wire bundle including neutral fraction was already shown in Fig. 12. A set of neutral fraction curves for a large range of current densities is shown in Fig. 13. As seen from these curves, there is always a range of angles, just outside of the beam before the signal strength declines into the noise level, in which the apparent neutral fraction is relatively constant. Observe also that the range changes with the current density. At higher current densities the beam expands more, and the neutral eye must be located at a large angle or made insensitive to the presence of the energetic ions. The 15° position of the neutral detector appears to be adequate for current densities of up to about 15 ma/cm<sup>2</sup>.

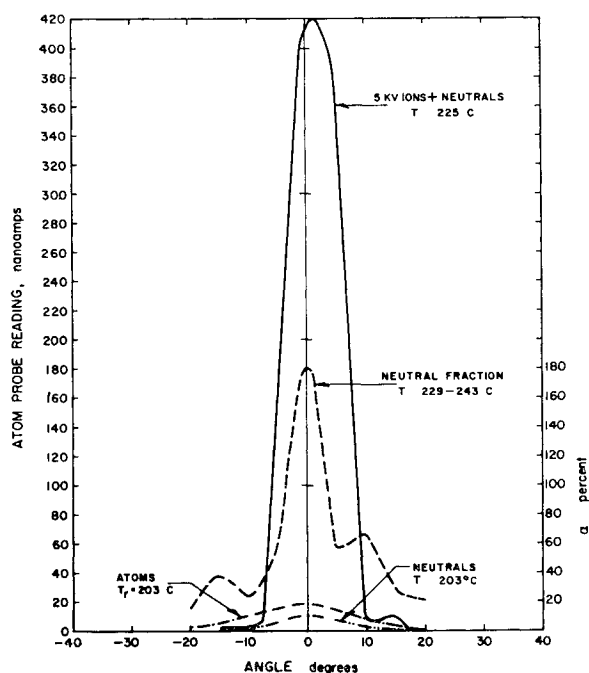
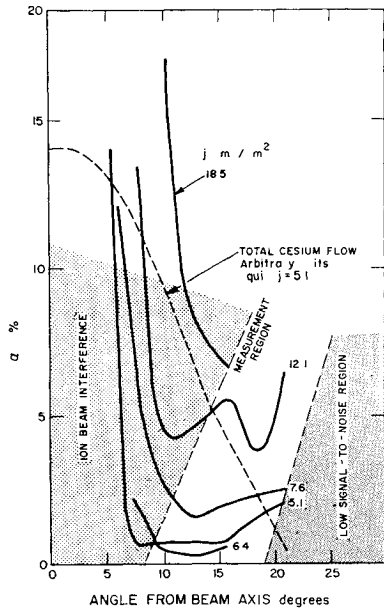
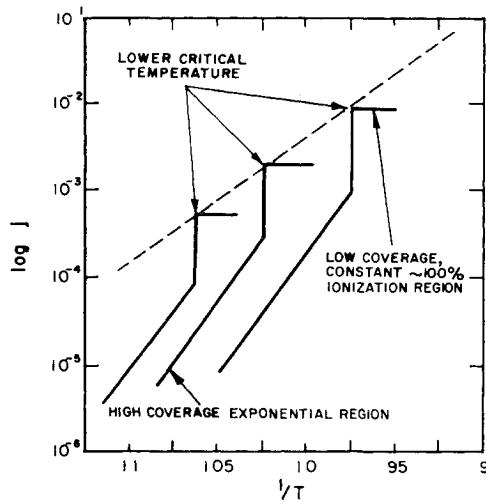


Fig. 12 Angular distribution of ions, atoms, neutrals, and neutral fraction through electrodes, WB5



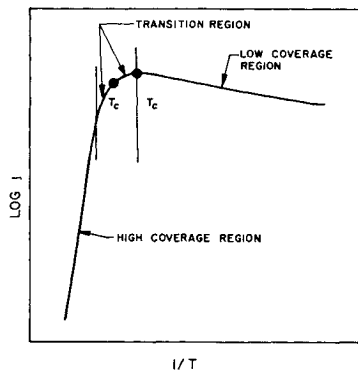
**Fig 13** Angular dependence of apparent neutral fraction and total cesium flow (as neutrals) for spherical powder ionizer, SP2



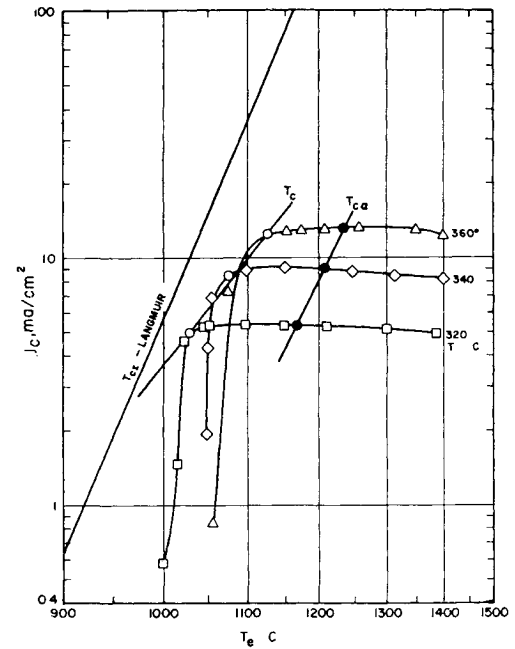
**Fig 14** Definitions of surface ionization regions for ion current-ionizer temperature curves for solid tungsten

## 9 Ion Current Density, Neutral Fraction, and Critical Temperatures

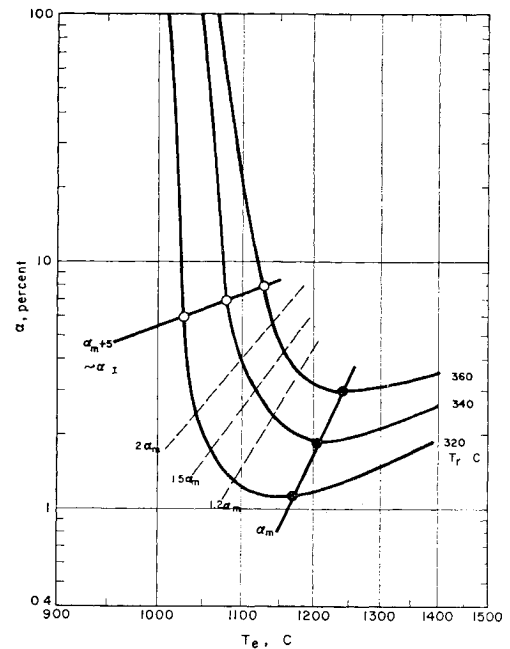
The most popular measurement in surface ionization is the ion current density as a function of ionizer temperature for various cesium flow rates. A set of typical curves for solid tungsten as measured by Langmuir<sup>11</sup> is shown in Fig 14 to define the various regions of operation. These data are usually taken with decreasing temperatures. At high temperature the cesium coverage is low, and for high work function surfaces the ionization is so nearly 100% that it appears essentially constant. At a certain temperature, called the



**Fig 15** Definitions of surface ionization regions for cesium on typical porous tungsten ionizers



**Fig 16a** Log  $j$  vs  $1/T$  curves defining two "critical temperature" points for SP1 ionizer



**Fig 16b** Log  $\alpha$  vs  $1/T$  curves showing neutral fraction corresponding to curves in Fig 16a

critical temperature, the ion current discontinuously drops by more than a decade as the surface becomes covered with cesium. In the high-coverage region the decline is exponential.

For the porous ionizer, the typical current-temperature curve is more like that in Fig 15. The low-coverage region is most likely no longer constant because of a change in ionization efficiency and ionizer permeability with ionizer temperature. Instead of a discontinuity, there is a transition region characterized by a large change in the slope. There are sometimes two "critical temperatures": one,  $T_{c1}$ , defined as the temperature at which the ion current has declined by 5%, and the other,  $T_c$ , defined as the temperature corresponding to minimum neutral fraction.

A typical set of selected data for SP1 ionizer illustrates the various definitions and graphing techniques in Fig 16. The



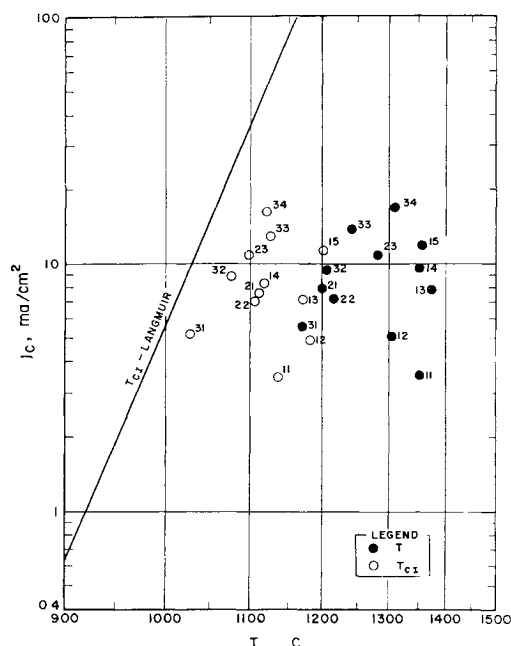


Fig 17a Log  $j$  vs  $1/T_c$ : summary of all critical points for SP1 ionizer

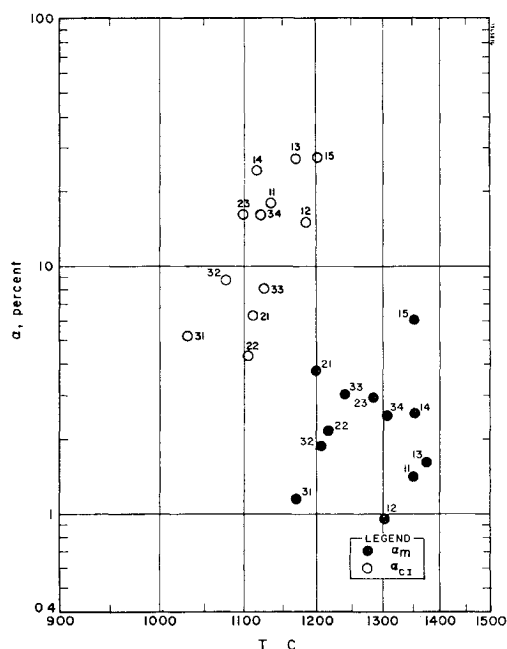


Fig 17b Log  $\alpha$  vs  $1/T$ : summary of all critical points for SP1 ionizer

current density  $j$  is based upon an area of  $0.179 \text{ cm}^2$  and was not corrected for 12% braze closure. Hence, the true ion current density is 12% higher. All curves were taken with decreasing ionizer temperatures, and the  $T_{C1}$  points are the lower critical temperature, which may be below the minimum operating temperature for the ionizer. The transition with increasing temperatures can occur at a higher temperature (about  $10\text{--}30^\circ$ ) called the upper critical temperature. The log  $j$  vs  $1/T$  plots are used because, for solid tungsten, the  $T_{C1}$  points fall on a line marked  $T_{C1}$ -Langmuir. The  $T$  parameter indicated is the temperature of the cesium reservoir in degrees Centigrade.

Simultaneous with the ion current, the neutral fraction is measured, and the corresponding points are plotted in Figs 16a and 16b. Sometimes the temperature  $T_{C\alpha}$  corresponding

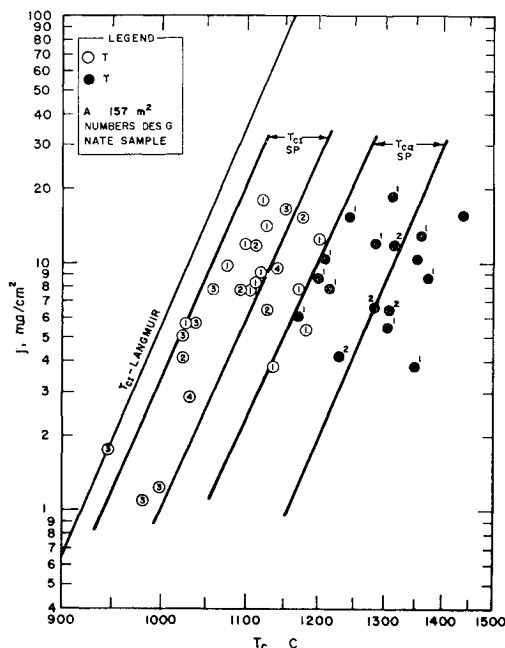


Fig 18 Summary of current and neutral fraction critical temperature points for four spherical powder ionizers

to a minimum neutral fraction of the SP1 ionizer should be operated at least  $100^\circ\text{C}$  hotter than the  $T_{C1}$  line. The  $T_{C1}$  points in the  $j$  vs  $1/T$  plot coincide closely to the  $\alpha_m = +5\%$  line in the  $\alpha$  vs  $1/T$  plot. If 6–8% neutral fractions cannot be tolerated, the ionizer must be operated at a temperature above the  $T_{C1}$  line. However, one does not have to go to  $\alpha_m$ . As the lines  $1.5\alpha_m$ ,  $1.5\alpha_m$ , and  $2\alpha_m$  indicate, the change of  $\alpha$  near the minimum is very slow. A  $30^\circ$  increase in temperature decreases  $\alpha$  by more than 50%. Optimum ionizer operation is, therefore, close to the  $T_{C1}$  line and between  $T_{C1}$  and  $T_{C\alpha}$  in Fig 16a, corresponding to about the  $1.5\alpha_m$  line in Fig 16b.

Figures 17a and 17b contain a summary of all critical temperature points for SP1 ionizer obtained from points on all curves similar to those in Figs 16a and 16b. The number with each point indicates the day and order in which the curves were taken. For example, 32 was taken during the third day of testing, and it is the second curve. Such graphical summaries of critical points for each ionizer with the order identified show the time trends, especially during the initial period of ionizer conditioning. As seen by comparison with Fig 16a, the critical temperature  $T_{C1}$  and  $T_{C\alpha}$  has decreased by about  $100^\circ$  from the first day to the third day. Figure 17b summarizes all neutral fraction values  $\alpha_{C1}$  corresponding to  $T_{C\alpha}$ , showing the time measurements. The  $\alpha_m$  values for this ionizer are much lower than the  $\alpha_{C1}$ . The  $\alpha_{C1}$  is very nearly equal to  $\alpha_m + 5$ . For this ionizer, curves on the first day showed no definite minimum up to  $1300^\circ\text{C}$ , and their slopes were lower.

To establish the performance of a given porous structure, several ionizers of the same structure must be tested. The current and neutral fraction vs critical temperature points of four spherical powder ionizers are shown in Fig 18. Neglecting the points measured during the conditioning period, the current vs critical temperature points are in a band marked  $T_{C1}$ -SP. This band is about  $75^\circ\text{C}$  wide and starts about  $25^\circ\text{C}$  above the Langmuir line. The neutral fraction critical temperature points fall in a similar band marked  $T_{C\alpha}$ -SP which is  $125^\circ\text{C}$  wide and starts about  $150^\circ\text{C}$  above the Langmuir line.

A similar summary of  $T_{C1}$  points for commercial sintered tungsten ionizers, given in Fig 19, shows a larger spread. For most SW ionizers, the  $T_{C\alpha}$  points were identical with  $T_{C1}$ .

A typical set of ion current density and neutral fraction curves for a  $12\text{-}\mu$  wire bundle ionizer with critical temperature



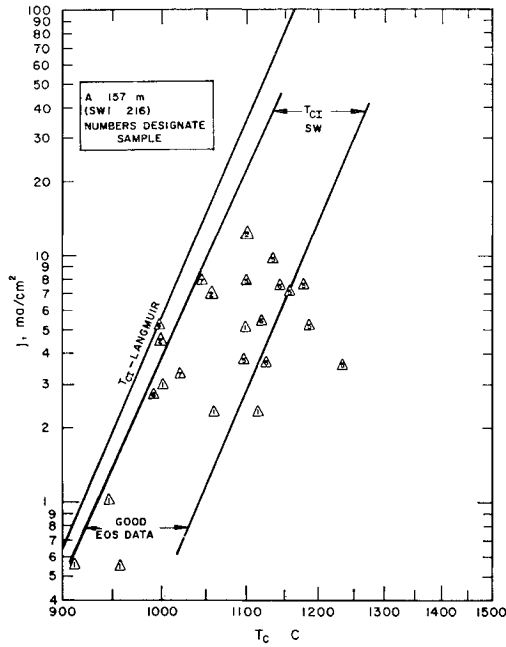


Fig 19 Summary of  $T_{CI}$  points for five commercial sintered tungsten ionizers

points is shown in Fig 20. Note the much higher neutral fractions and the coincidence of most  $T_{CI}$  and  $T_{C\alpha}$  points except at low  $\alpha$  values.

The fine wire bundle ionizer  $\alpha$  curves in Fig 21 showed much more pronounced minima, especially at low flow rates. Because of ionizer-to-reservoir heat flow, the reservoir temperature depended upon the ionizer temperature. The nature of the  $\alpha$  curves changed with ionizer conditioning. Numbers with each curve again indicate their order.

The  $\log \alpha$  vs  $T_i$  graphs (Fig 16b) are most useful in establishing the neutral fraction critical temperature  $T_{C\alpha}$  and the values of  $\alpha_{CI}$  at the ionizer critical temperature. To compare the neutral fractions of various ionizers and structures, a plot of  $\alpha$  vs  $j$  is used. Figure 22 shows the minimum neutral fraction of several ionizers for the various porous tungsten structures. Some of the spread for each ionizer is caused by

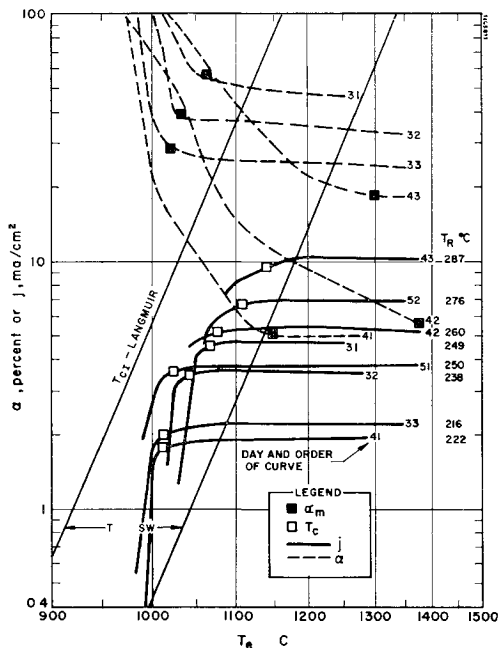


Fig 20 Ion current and neutral fraction vs ionizer temperature for coarse wire bundle WB3

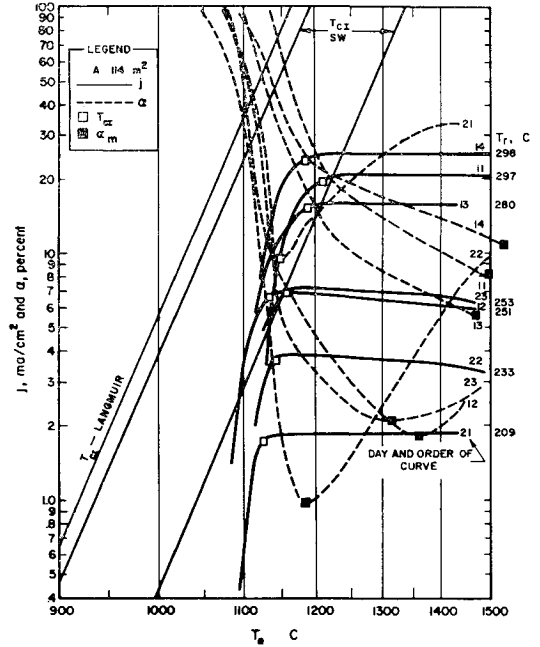


Fig 21 Ion current and neutral fraction vs ionizer temperature for fine wire bundle WB6

the time dependence illustrated in Fig 23. Note the large reduction during the initial conditioning and the smaller change on each successive day. The commercial porous tungsten ionizers, except for SW5, and the spherical powder ionizers generally have neutral fractions ranging from a fraction of a percent to 10% at 30 ma/cm<sup>2</sup>. This is higher than given by the Langmuir data for solid tungsten. The large neutral fractions of the 12- $\mu$  wire bundle ionizers are probably due to the opening of the pores at high temperatures and the appearance of small surface cracks. The fine wire bundles made of 6- $\mu$ -diam wire had much lower neutral fractions than the coarse wire bundles, after an initial conditioning period of about 1-2 hr.

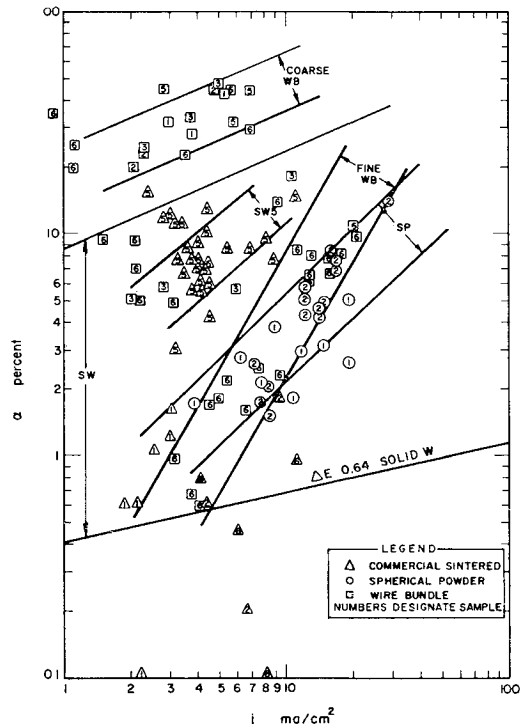


Fig 22 Neutral fractions for various porous tungsten structures

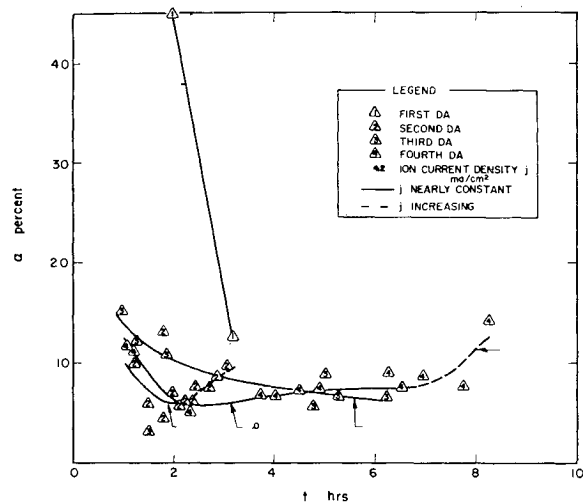
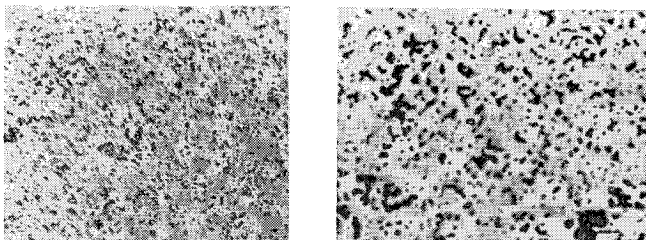


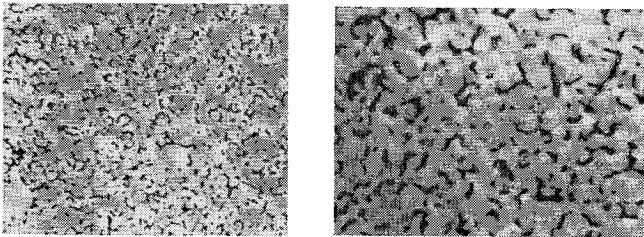
Fig 23 Time dependence of neutral fraction SW5

10 Comparison of Commercial, Spherical Powder, and Wire Bundle Ionizers

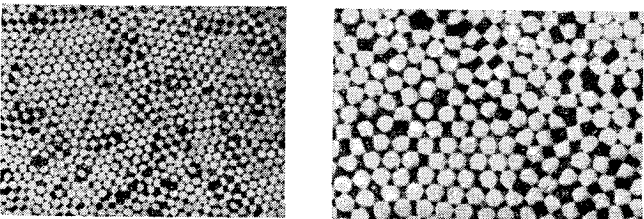
In Fig 24, surface photomicrographs show the pore sizes and distances for typical ionizers of each type. The commercial porous tungsten has irregular pore sizes and shapes with an equivalent diameter of  $2.4\ \mu$ , and an average pore distance of  $6.7\ \mu$ . The spherical powder ionizers have more regular slitlike pores with areas for which the equivalent diameter is  $6.4\ \mu$ . The average pore distance, assuming a regular distribution into a hexagonal close-packed array, is  $13.5\ \mu$ . The coarse wire bundle has some very large pores of about  $15\ \mu$  with an average of about  $4\ \mu$ . For most coarse WB ionizers made of 0.0005-in.-diam wire, the wire ends were expanded mechanically at the surface to reduce the pore size and produce uniform packing. Operation at high temperature tended



Commercial porous tungsten



Spherical powder



Wire bundle

Fig 24 Set of comparative surface photomicrographs

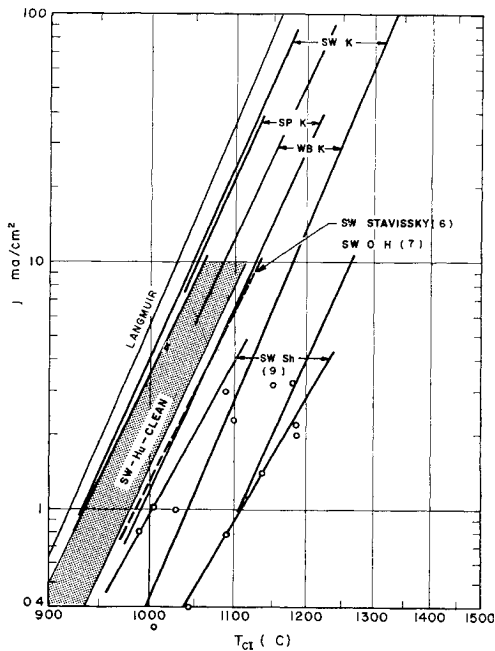


Fig 25 Comparison of critical temperatures for 1) commercial ionizers, 2) spherical powder, and 3) wire bundle ionizers

to open up the pores. The fine wire bundle made of 0.00025-in.-diam wire had a 0.3- $\mu$  equivalent pore diameter.

The apparent nitrogen permeability of all ionizers is compared in Table 2. The average SP ionizer was about 4 times the average WB ionizer, and about 10 times more permeable to nitrogen than the average SW ionizer. The cesium conductance at operating temperature as given by ion current density in Fig 6 follows a similar pattern, but the spread is larger for the atom detector data and smaller for the ion current density data. The change of permeability with the ionizer temperature was about 3%/100°K for all structures.

The angular distributions of atoms indicated evaporation from the surface for SW ionizers and from the surface for WB ionizers with small pore sizes or from within the pores for WB ionizers with large pore sizes.

Table 2 Apparent nitrogen permeability of mounted ionizers

Ionizer identity	$K, 10^{-6} \text{g}/(\text{cm torr-sec})$
Commercial sintered	
SW1	0.078
SW2	0.122
SW3	0.226
SW4	0.186
SW5	0.257
SW6	0.285
SW7	0.570
Average	0.246
Spherical powder	
SP1	0.37
SP2	0.493
SP3	0.628
SP4	0.553
SP5	1.21
Average	0.650
Wire bundle, 12 $\mu$	
WB1	1.41
WB2	1.37
WB3	2.90
WB4	3.34
WB5	
Average	2.25
Wire bundle, 6 $\mu$	
WB6	2.58

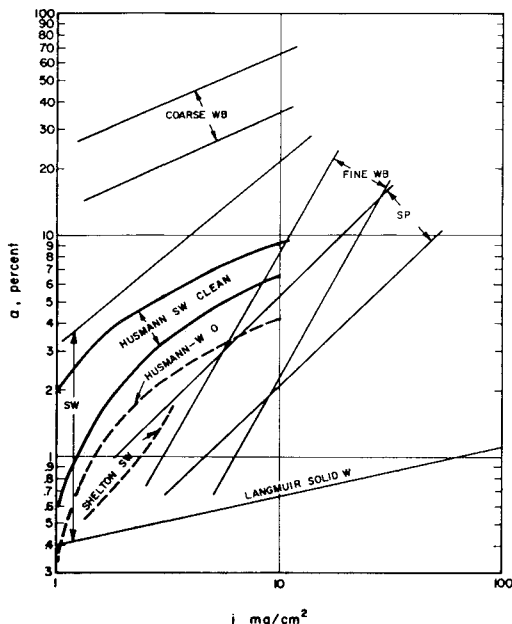


Fig 26 Comparison of neutral fractions of the various tungsten structures with results of other investigations

The critical temperatures of all *SP* ionizers given in Fig 18 demonstrate the spread of points from sample to sample and the difference between  $T_{CI}$  and  $T_{Ca}$ . Such differences were not observed for the *SW* ionizers in Fig 19 or for most of the 12- $\mu$  wire bundles as in Fig 20. A very pronounced minimum was observed in  $\alpha$  at  $T_{Ca}$ , which was very nearly equal to  $T_{CI}$  for the 6- $\mu$  wire bundle in Fig 21.

A comparison of current vs critical temperature  $T_{CI}$  data for all tungsten structures with data published by other investigators is given in Fig 25. Nearly all tungsten ionizers fall within a band that is between about 30° and 130°C above the Langmuir line. The data taken during initial conditioning periods are neglected. The spread for *SP* and *WB* ionizers is smaller (about 90°C) as compared to the *SW* ionizers.

Similar distributions of points are found in the  $\alpha$  vs  $j$  data of Fig 21. These same data are summarized and compared with results of other investigators in Fig 26. The coarse wire bundle ionizers had very high neutral fractions, and the commercial porous tungsten showed a very large spread. The spherical powder ionizers had a narrow spread with values similar to the fine wire bundle. The highest ion current density was reached with the *SP* ionizers, the fine wire bundle being next highest. Again all data, except the band due to Hussman and marked *SW-clean*, were obtained at  $10^{-6}$ -torr

residual gas pressure where the surfaces are probably oxygenated.

Performance tests are continuing, and some ionizers of each type will be tested at lower residual gas pressures in an ultra-high vacuum system. The more regular structures, such as the fine wire bundle and the spherical powder ionizers, have given better performance than the commercial sintered tungsten, but a final comparison should include data on more samples and at lower residual gas pressures.

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